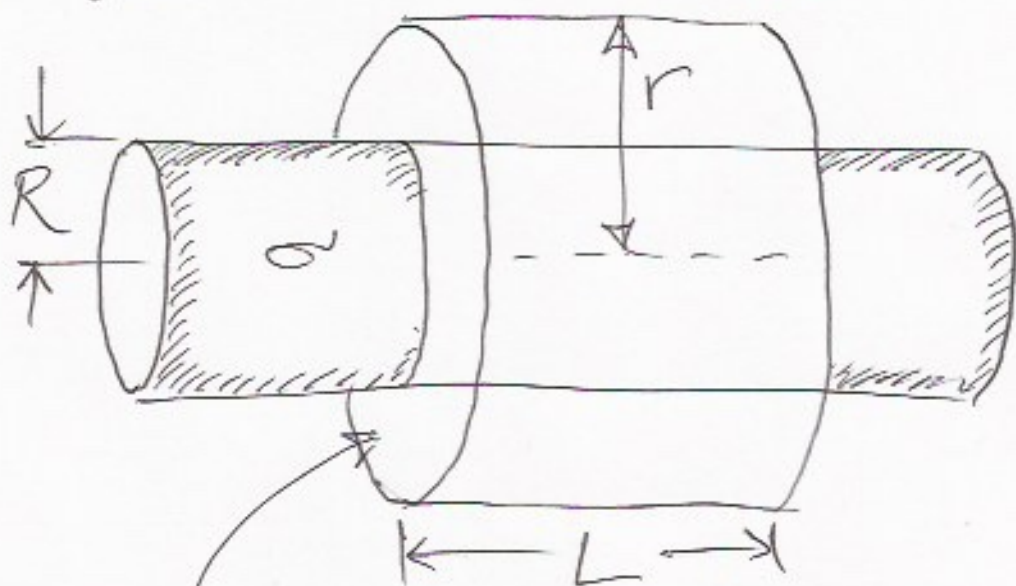


Assignment 9 Solutions

(1)



$S(r, L) = \text{cylinder surface}$

By symmetry: $\vec{E} = E(r) \hat{r}$

$$\text{Gauss: } \oint_{S(r, L)} \vec{E} \cdot d\vec{a} = E(r) \cdot 2\pi r L$$

$$= \frac{Q_{\text{enc.}}}{\epsilon_0} = \begin{cases} 0, & r < R \\ \frac{2\pi r L \sigma}{\epsilon_0}, & r > R. \end{cases}$$

$$\Rightarrow E(r) = \begin{cases} 0, & r < R \\ \frac{\sigma}{\epsilon_0} \frac{R}{r}, & r > R. \end{cases}$$

Since for $r > R$ the electric field has exactly the same form as that of an infinite line of charge, we can use the results from lecture in that region: (2)

$$\lambda \rightarrow 2\pi R\epsilon$$

$$\vec{F} = \begin{cases} 0, & r < R \\ -q\epsilon B \hat{r}, & r > R \end{cases}$$

$$B = \frac{\mu_0 I}{2\pi r}, \quad I = 2u\lambda = 4\pi u R \epsilon$$

1. The "double-curl" identity is an instance of the double-cross-product identity, where two of the vectors are the gradient operator:

$$\vec{\nabla} = \frac{\partial}{\partial x} \hat{x} + \frac{\partial}{\partial y} \hat{y} + \frac{\partial}{\partial z} \hat{z}$$

$$\vec{A} = \vec{\nabla}, \quad \vec{B} = \vec{\nabla}, \quad \vec{C} = \vec{V} \quad (3)$$

$$(*) \quad \vec{A} \times (\vec{B} \times \vec{C}) = \vec{B}(\vec{A} \cdot \vec{C}) - (\vec{A} \cdot \vec{B})\vec{C}$$

Note that the right-hand-side has \vec{C} always on the right, so the derivatives in \vec{A} & \vec{B} will act on it (\vec{V}).

To check (*), just work out the z-component on both sides, since the other two components are related by symmetry:

$$(A_x \hat{x} + A_y \hat{y} + \dots) \times ((B_y C_z - B_z C_y) \hat{x} + (B_z C_x - B_x C_z) \hat{y} + \dots)$$

$$= \left(\overset{(1)}{A_x} (\overset{(2)}{B_z C_x - B_x C_z}) - \overset{(3)}{A_y} (\overset{(4)}{B_y C_z - B_z C_y}) \right) \hat{z} + \dots$$

$$\begin{aligned}
 & \vec{B}(\vec{A} \cdot \vec{C}) - (\vec{A} \cdot \vec{B})\vec{C} \quad (4) \\
 & = \left(B_z \left(\overset{(1)}{A_x C_x} + \overset{(4)}{A_y C_y} + \cancel{A_z C_z} \right) \right. \left. \begin{array}{l} \uparrow \text{cancel} \\ \downarrow \end{array} \right. \\
 & \quad \left. - \left(\overset{(2)}{A_x B_x} + \overset{(3)}{A_y B_y} + \cancel{A_z B_z} \right) C_z \right) \hat{z} \\
 & \quad \quad \quad + \dots
 \end{aligned}$$

The four terms that do not cancel exactly match the four terms on the previous page.

2. Let Φ_0 be the point-charge solution:

$$\Phi_0(\vec{r}; \vec{r}') = \frac{q}{4\pi\epsilon_0 |\vec{r} - \vec{r}'|}$$

\vec{r} = field point, \vec{r}' = source point

Claimed general solution, for charge density ρ , is (next page)

$$\Phi(\vec{r}) = \int d^3r' \frac{\rho(\vec{r}')}{\epsilon_0} \Phi_0(\vec{r}; \vec{r}') \quad (5)$$

To check, first apply ∇_r^2 to both sides (derivatives act only on field point):

$$\begin{aligned} \nabla^2 \Phi(\vec{r}) &= \int d^3r' \frac{\rho(\vec{r}')}{\epsilon_0} \underbrace{\nabla_r^2 \Phi_0(\vec{r}; \vec{r}')}_{-\frac{\delta^3(\vec{r}-\vec{r}')}{\epsilon_0}} \\ &= -\frac{\rho(\vec{r})}{\epsilon_0} \quad \checkmark \end{aligned}$$

Next, note that the Φ defined by the integral has $\Phi=0$ boundary condition at ~~the~~ $|\vec{r}| \rightarrow \infty$, since

$\Phi_0(\vec{r}; \vec{r}') \rightarrow 0$, $|\vec{r}| \rightarrow \infty$ for any finite $|\vec{r}'|$.

It's not possible to have two (6)
(different) solutions of the Poisson
equation for the same ρ and zero
boundary condition at $|r| = \infty$,

$$-\nabla^2 \phi_1 = \rho/\epsilon_0, \quad -\nabla^2 \phi_2 = \rho/\epsilon_0,$$

since then $-\nabla^2(\phi_1 - \phi_2) = 0$ would
be another solution ^{ϕ_3} of Laplace
equation, with zero boundary condition
at infinity, in addition to $\phi_3 = 0$.
(Solutions to the Laplace equation
are unique.)

3. We are given \vec{V} and need to
find f such that

$$\vec{V}' = \vec{V} + \vec{\nabla} f \quad \text{has zero divergence.}$$

$$0 = \vec{\nabla} \cdot \vec{V}' = \vec{\nabla} \cdot \vec{V} + \nabla^2 f$$

(7)

$$\Rightarrow -\nabla^2 f = \vec{\nabla} \cdot \vec{V} = \rho / \epsilon_0$$

(It needn't have the interpretation of a charge density, but we can use the same symbol.)

$$\Rightarrow f(\vec{r}) = \int d^3r' \frac{\vec{\nabla} \cdot \vec{V}'(\vec{r}')}{4\pi |\vec{r} - \vec{r}'|}$$

Since $\vec{\nabla} \times \vec{V}' = \vec{\nabla} \times \vec{V} + \underbrace{\vec{\nabla} \times \vec{\nabla} f}_0$,

this shows it's always possible to modify a vector field,

$$\vec{V} \rightarrow \vec{V}'$$

without changing its curl, so that it has zero divergence.