

Assignment 8 Solutions

(1)

electric field flux for moving point charge :

$S(r)$ = surface of sphere, radius r

$$\Phi = \oint_{S(r)} \vec{E} \cdot d\vec{a} = \int_{\downarrow 2\pi} r^2 d\phi \sin\theta d\theta \frac{Kq}{r^2} \frac{1}{\gamma^2} \frac{1}{(1-\beta^2 \sin^2\theta)^{3/2}}$$

$$= 2\pi Kq \int_0^\pi \frac{\sin\theta d\theta}{\gamma^2 (1-\beta^2 \sin^2\theta)^{3/2}}$$

$$2\pi K = \frac{1}{2\epsilon_0}, \quad \sin\theta d\theta = -\underbrace{d(\cos\theta)}_u$$

$$1-\beta^2 \sin^2\theta = 1-\beta^2 + \beta^2 u^2$$

$$\frac{1}{\gamma^2} = 1-\beta^2$$

$$\Phi = \frac{q}{2\epsilon_0} \int_{-1}^1 \frac{(1-\beta^2)(+du)}{((1-\beta^2) + \beta^2 u^2)^{3/2}}$$

useful "derivative fact":

(2)

$$\begin{aligned}\frac{d}{du} \left(\frac{u}{\sqrt{A+Bu^2}} \right) &= \frac{1}{\sqrt{A+Bu^2}} - \frac{1}{2} \frac{2Bu^2}{(A+Bu^2)^{3/2}} \\ &= \frac{A+Bu^2 - Bu^2}{(A+Bu^2)^{3/2}} \\ &= \frac{A}{(A+Bu^2)^{3/2}}\end{aligned}$$

set $A = 1 - \beta^2$, $B = \beta^2$

$$\underline{\Phi} = \frac{q}{2\epsilon_0} \int_{-1}^1 \frac{A du}{(A+Bu^2)^{3/2}} = \frac{q}{2\epsilon_0} \left. \frac{u}{\sqrt{A+Bu^2}} \right|_{-1}^1$$

$$= \frac{q}{2\epsilon_0} \left(\frac{1}{\sqrt{A+B}} - \frac{-1}{\sqrt{A+B}} \right) = \frac{q}{2\epsilon_0} \times 2$$

$$= \frac{q}{\epsilon_0} \checkmark$$

Field-line pattern of oscillating point charge: (3)

At $t=0$ the charge is at $x=0$ and moving along the x -axis with velocity $+c/2$.

Define the following times & positions:

t_a : times in the past when the charge accelerated (reversed its velocity)

x_a : positions ~~of~~ of charge at the times t_a

x_+ : position of light pulse ~~moving~~ moving with velocity $+c$ and launched at x_a & t_a , now at $t=0$.

x_- : same as x_+ , but for light pulse moving with velocity $-c$.

$$x_+ = x_a + c(t - t_a), \quad x_- = x_a - c(t - t_a) \quad (4)$$

x_0 : extrapolated position of charge if it didn't undergo accelerations after t_a

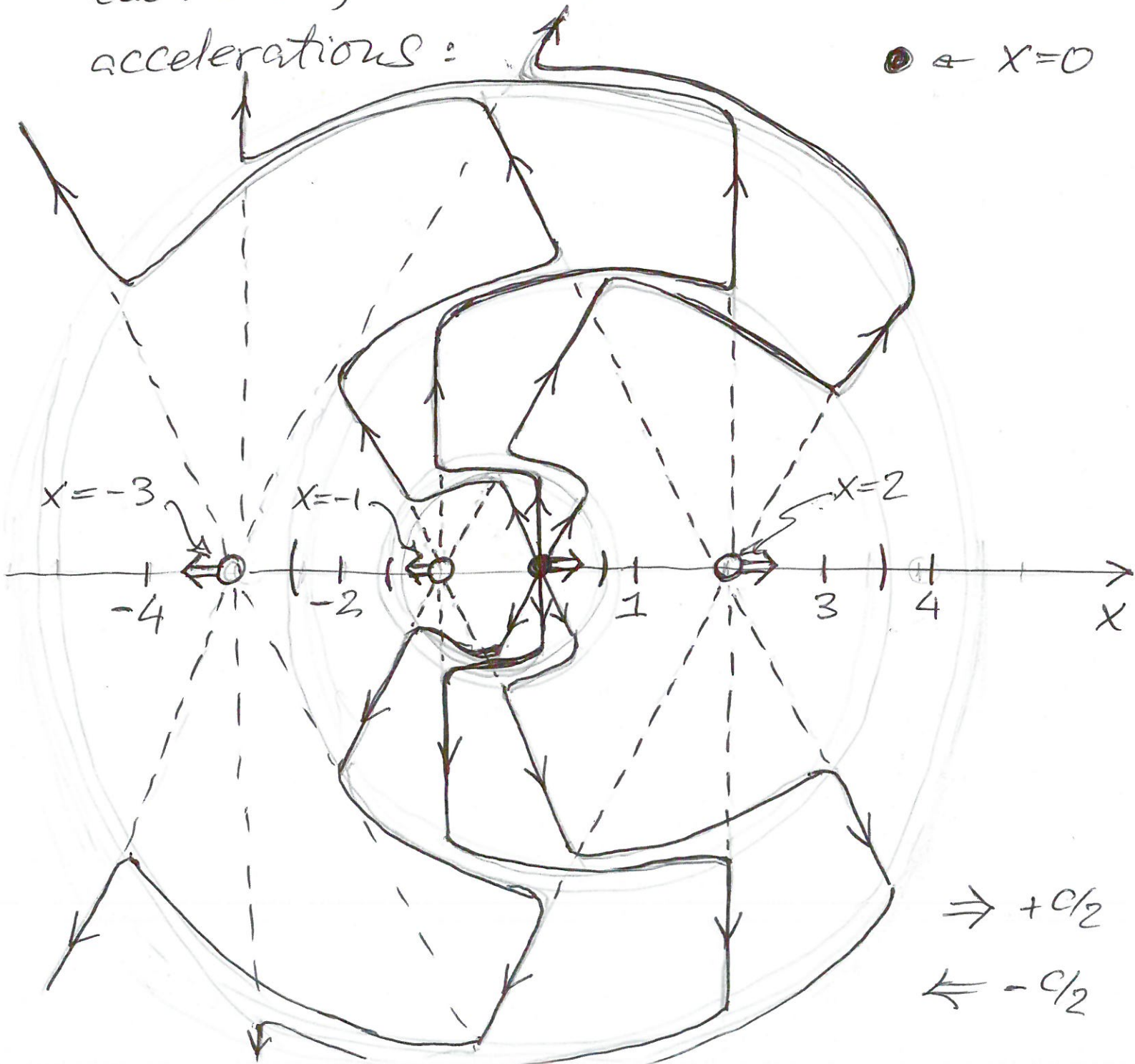
$$x_0 = x_a + \frac{c}{2}(t - t_a) \quad \text{if moving with pos. velocity}$$

$$x_0 = x_a - \frac{c}{2}(t - t_a) \quad \text{if neg. velocity}$$

Here is a table of these times/positions for $t=0$. Times are in units of ns, positions in ft., so $c \cong 1$.

t_a	-1	-3	-5	-7
x_a	-1/2	+1/2	-1/2	+1/2
x_+	1/2	7/2	9/2	15/2
x_-	-3/2	-5/2	-11/2	-13/2
x_0	0	-1	+2	-3

The positions $x_+ \approx x_-$ are the 5
 intersection of each "light-pulse shell"
 with the x -axis and x_0 is where
 the field lines would converge inside
 each shell, in the absence of later
 accelerations:



charge in uniform electric field (6)

$$\frac{d\vec{P}}{dt} = qE\hat{x}, \quad \vec{P}(0) = 0 \quad (\text{at rest when } t=0)$$

$$P_y = P_z = 0$$

$$\underline{P_x = qEt}$$

$$P_x = \frac{m\gamma v_x}{(1 - v_x^2/c^2)^{1/2}}$$

solve for v_x
in terms of P_x

$$v_x = \frac{P_x/m}{(1 + (P_x/mc)^2)^{1/2}}$$

$$= \frac{(\frac{qE}{m})t}{(1 + (\frac{qE}{mc})^2 t^2)^{1/2}}$$

define a "time" $\tau = \frac{mc}{qE}$

$$v_x = \frac{c(t/\tau)}{(1 + (t/\tau)^2)^{1/2}}$$

From the last expression we see that for $t \rightarrow +\infty$, $v_x \rightarrow c$, and v_x is already close to c when $t \gg \tau$. In the infinite past, or $t \rightarrow -\infty$, $v_x \rightarrow -c$. In other words, the charge had been moving with negative, near-light velocity before the electric field brought it to rest, instantaneously, at $x=0$ and time $t=0$.

$$\begin{aligned}
 x &= \int_0^t v_x dt = c \int_0^t \frac{(t/\tau) dt}{(1 + (t/\tau)^2)^{1/2}} \\
 &= c\tau \sqrt{1 + \frac{t^2}{\tau^2}} \Big|_0^t = c\tau \left(\sqrt{1 + \frac{t^2}{\tau^2}} - 1 \right)
 \end{aligned}$$