

Assignment 5 Solutions

Easy evaluation of the electric energy density integral

Prove the identity

$$\nabla \cdot (\mathbf{E}_1 \varphi_2) = (\nabla \cdot \mathbf{E}_1) \varphi_2 + \mathbf{E}_1 \cdot \nabla \varphi_2$$

Each of the three terms is itself the sum of three terms, one for each axis in three dimensions. Here are just the terms involving $\partial/\partial x$ and x -components of vectors:

$$\frac{\partial}{\partial x} (E_{1x} \varphi_2) = \frac{\partial E_{1x}}{\partial x} \varphi_2 + E_{1x} \frac{\partial \varphi_2}{\partial x}.$$

But this is just the ordinary product rule of calculus. The same identity applies to the y and z counterparts. Summing the three product-rule identities gives the identity you were asked to prove.

Use the divergence theorem to prove that

$$\int d^3 \mathbf{r} \nabla \cdot (\mathbf{E}_1 \varphi_2) = 0.$$

By the divergence theorem,

$$\int_{\mathcal{V}} d^3 \mathbf{r} \nabla \cdot (\mathbf{E}_1 \varphi_2) = \oint_S (\mathbf{E}_1 \varphi_2) \cdot d\mathbf{a},$$

where S is the closed surface that bounds the region \mathcal{V} . We are free to take \mathcal{V} to be a large sphere of radius R that encloses both point charges with center fixed relative to these charges. In the limit of large R the difference between the sphere center and the positions of the charges is unimportant and $|\mathbf{E}_1| \sim c/R^2$ and $\varphi_2 \sim c'/R$, where c and c' are constants, so that the integrand of the surface integral behaves as $1/R^3$ for large R . Since the surface area of the sphere only grows as R^2 , the surface integral decays as $1/R^3 \times R^2 = 1/R$ and is zero in the limit of an infinite sphere (corresponding to the integral over \mathcal{V} being the integral over all space).

You are now left with evaluating the integral

$$\int d^3 \mathbf{r} (\nabla \cdot \mathbf{E}_1) \varphi_2.$$

You can evaluate this without doing any serious work. The divergence of the electric field produced by charge 1 evaluates to a simple distribution; express it in terms of the Dirac delta “function”.

The source of the field \mathbf{E}_1 is a point charge q_1 located at \mathbf{r}_1 , that is, the charge density

$$\rho_1(\mathbf{r}) = q_1 \delta^3(\mathbf{r} - \mathbf{r}_1) .$$

Since $\nabla \cdot \mathbf{E}_1 = \rho_1/\epsilon_0$ (local Gauss law), we can rewrite the integral as

$$\int d^3\mathbf{r} (\rho_1(\mathbf{r})/\epsilon_0)\varphi_2(\mathbf{r}) = \int d^3\mathbf{r} q_1 \delta^3(\mathbf{r} - \mathbf{r}_1)\varphi_2(\mathbf{r})/\epsilon_0 .$$

Using the delta-function identity this evaluates to

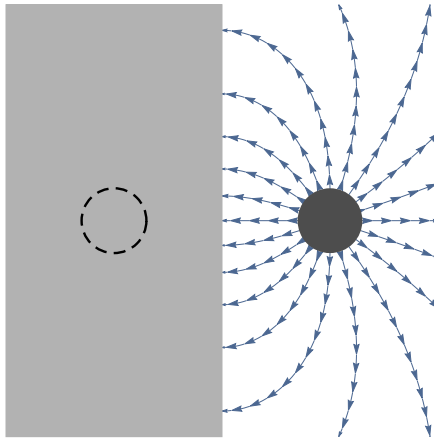
$$(q_1/\epsilon_0)\varphi_2(\mathbf{r}_1) = \frac{K}{\epsilon_0} \frac{q_1 q_2}{|\mathbf{r}_1 - \mathbf{r}_2|} .$$

Your answer for U_{12} should equal the answer you got in assignment 2 the hard way.

$$\begin{aligned} U_{12} &= \epsilon_0 \int d^3\mathbf{r} \mathbf{E}_1(\mathbf{r}) \cdot \mathbf{E}_2(\mathbf{r}) \\ &= -\epsilon_0 \int d^3\mathbf{r} \mathbf{E}_1(\mathbf{r}) \cdot \nabla \varphi_2(\mathbf{r}) \\ &= \epsilon_0 \int d^3\mathbf{r} (\nabla \cdot \mathbf{E}_1)\varphi_2 \\ &= K \frac{q_1 q_2}{|\mathbf{r}_1 - \mathbf{r}_2|} . \end{aligned}$$

Point charge near the surface of a conductor

Keeping in mind the rules that apply to field lines near a conductor, make a drawing of the field lines in the x - y plane (the 3D pattern is symmetric about the x -axis). If this reminds you of a field line drawing you've seen before, state which it was!



This resembles the field of a dipole in the region $x > 0$. If we had a true dipole, there would also be a charge $-q$ at the center of the dashed circle. By symmetry, the dipole field lines are perpendicular to the plane $x = 0$, exactly as required at the surface of a conductor (the region $x < 0$). In the case of a conductor, instead of a charge $-q$, the field lines terminate at the conductor because there is a negative charge density on its surface.

Compute the magnitude of this surface field, $E(r)$, as a function of the distance r from the x -axis. As a check, your answer should have the form

$$E(r) = \frac{A}{(r^2 + B)^{3/2}},$$

where A and B are constants (to be determined by you).

The electric field produced by the two point charges, one real and the other fictitious, evaluated at $x = 0$ (technically, just outside the conductor), and arbitrary y is

$$\begin{aligned} \mathbf{E} &= Kq \left(\frac{-a\hat{\mathbf{x}} + y\hat{\mathbf{y}}}{(a^2 + y^2)^{3/2}} - \frac{a\hat{\mathbf{x}} + y\hat{\mathbf{y}}}{(a^2 + y^2)^{3/2}} \right) \\ &= Kq \frac{-2a\hat{\mathbf{x}}}{(a^2 + y^2)^{3/2}}. \end{aligned}$$

This is perpendicular to the surface of the conductor (the $x = 0$ plane) and has magnitude of the form given in the assignment with

$$A = 2Kqa, \quad B = a^2.$$

Calculate $\sigma(r)$, using Gauss's law, as we did in lecture. Integrate this surface charge density over the entire plane $x = 0$ to find out how much total charge resides on the surface of the conductor.

As shown in lecture, the electric field magnitude at the surface of a conductor and the surface charge density there are related as $E = \sigma/\epsilon_0$. The sign of σ is positive when the field lines point away from the conductor, so in this case σ is negative and has value

$$\sigma(r) = -\epsilon_0 E(r) = -\frac{q}{2\pi} \frac{a}{(a^2 + r^2)^{3/2}}$$

since $\epsilon_0 K = 1/(4\pi)$. Integrating this density over the entire conductor surface,

$$\int_0^\infty (2\pi r dr) \sigma(r) = -q ,$$

we see that the system (point charge $+q$ and surface charge) is neutral.

Calculate the electric force, magnitude and direction, experienced by the $+q$ charge using whichever source (real surface charge or fictitious point charge) is simpler for the task.

From the local perspective it is the electric field \mathbf{E} at the location of the $+q$ charge that is responsible for the force $\mathbf{F} = q\mathbf{E}$ on the charge. We know that \mathbf{E} is produced by the surface charges, but exactly the same \mathbf{E} (and force) would be produced by a $-q$ charge at $x = -a$ and no surface charge:

$$\mathbf{F} = q \left(K \frac{(-q)}{(2a)^2} \hat{\mathbf{x}} \right) = -K \frac{q^2}{4a^2} \hat{\mathbf{x}} .$$