

## Coupled Pendulums

- The diagram on the right might be useful. The component of the force due to gravity perpendicular to the pendulum rod is  $-mg \sin(\theta_n)$ , and it's negative because it acts to decrease  $\theta_n$ . The torque due to gravity is then

$$\tau_{ng} = -mgr \sin(\theta_n)$$

- In class it was shown that the torque on pendulum  $n$  due to the neighboring pendulums was  $\tau_n = \kappa(\theta_{n-1} - 2\theta_n + \theta_{n+1})$ . The net torque on the  $n$ th pendulum is  $\tau_{n,net} = -mgr \sin(\theta_n) + \kappa(\theta_{n-1} - 2\theta_n + \theta_{n+1})$ . Applying Newton's second law for torques, we have that the equations of motion are

$$I \frac{d^2\theta_n}{dt^2} = -mgr \sin(\theta_n) + \kappa(\theta_{n-1} - 2\theta_n + \theta_{n+1})$$

Or equivalently,

$$\frac{d^2\theta_n}{dt^2} = -\frac{mgr}{I} \sin(\theta_n) + \frac{\kappa}{I} (\theta_{n-1} - 2\theta_n + \theta_{n+1})$$

Where  $I$  is the moment of inertia, and is a function of  $m$  and  $r$ .

- For the case when  $\theta_n$  is small, we have that  $\sin(\theta_n) \approx \theta_n$ . In this case, the equation becomes

$$\frac{d^2\theta_n}{dt^2} = -\frac{mgr}{I} \theta_n + \frac{\kappa}{I} (\theta_{n-1} - 2\theta_n + \theta_{n+1}).$$

- As shown in class, the second term can be interpreted in terms of the second derivative with respect to position, and the equation becomes

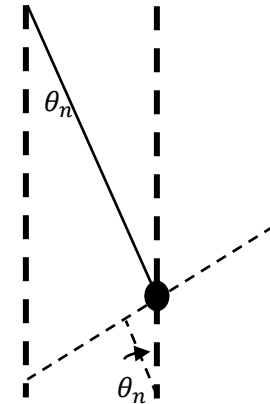
$$\frac{\partial^2 \theta(t, x)}{\partial t^2} = -\frac{mgr}{I} \theta(t, x) + \frac{\kappa \Delta x^2}{I} \frac{\partial^2 \theta(t, x)}{\partial x^2}.$$

## Co-tape™

- With the assumptions we are making (i.e. constant electric field, no edge effects) the tape and metal surfaces form a parallel plate capacitor. In such a capacitor, the field is given by  $E = \frac{\sigma}{\kappa \epsilon_0}$ , where  $\sigma$  is the charge per unit area. This is simply the sum of the fields of two planes of charge. Substituting in  $\sigma = \frac{Q}{Lw}$ , where  $l$  and  $w$  are the length and width of the tape, we have  $E = \frac{Q}{\kappa \epsilon_0 Lw}$ . The potential is then  $V = \int_0^d E(z) dz = \frac{Qd}{\kappa \epsilon_0 Lw}$ . With this we have that the capacitance  $C = \frac{Q}{V} = \frac{\kappa \epsilon_0 Lw}{d}$ . Finally, the capacitance per unit length is

$$\tilde{C} = \frac{\kappa \epsilon_0 w}{d}.$$

- Ignoring the edge effects, we compute the magnetic field as the sum of the magnetic fields of two sheets of current. The current density  $j = \frac{i}{w}$ , the current per unit width. For an arbitrary width  $\Delta x$ , using Ampere's law, we have,  $B_{sheet} 2\Delta x = \mu_0 j \Delta x = \frac{\mu_0 i \Delta x}{w}$ . The field



between the sheets is then  $B = \frac{\mu_0 i}{w}$ . The flux is  $\phi = \int_0^d \int_0^l B \, dy \, dz = \frac{\mu_0 i l d}{w}$ . The inductance is  $L = \frac{\phi}{i} = \frac{\mu_0 l d}{w}$ . The inductance per unit length is then

$$\tilde{L} = \frac{\mu_0 d}{w}.$$

3. The velocity of waves is

$$v = \frac{1}{\sqrt{\tilde{L}\tilde{C}}} = \frac{1}{\sqrt{\frac{\kappa\epsilon_0 w}{d} \frac{\mu_0 d}{w}}} = \frac{1}{\sqrt{\kappa}} \frac{1}{\sqrt{\epsilon_0\mu_0}} = \frac{c}{\sqrt{\kappa}}$$